# Bonding analysis and stability on alternant $B_{16} \mathbf{N}_{16}$ cage and its dimers 

Bing Yin • Guo Wang • Niya Sa • Yuanhe Huang

Received: 17 December 2007 / Accepted: 12 March 2008 /Published online: 21 May 2008
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#### Abstract

Bonding analysis is performed on alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage based on a combined study of DFT with NBO method. The main feature of such analysis is the separation of bonding structure into two components: $\sigma$ skeleton and $\pi$ bond system. Each component is further decomposed into contributions from various NBOs, thus we obtain the details of bonding interactions of every BN unit. Based on these results, relative stability of four covalent dimers of $\mathrm{B}_{16} \mathrm{~N}_{16}$ is predicted and this prediction is verified by DFT calculations. So the possibility of forecasting properties of oligomers just from analysis on monomer is highlighted in this way.


Keywords DFT• Bonding analysis • NBO •
Relative stability

## Introduction

Due to its role of isoelectronic species to carbon fullerene, stoichiomeric boron nitride cage $(\mathrm{BN})_{\mathrm{n}}$ have been investigated extensively both in experiments and calculations to exploit their electronic, optical and magnetic properties [115]. There are mainly two structural classes for $(\mathrm{BN})_{\mathrm{n}}$ cages [6, 12]: One is constructed from alternant BN units and governed by an isolated square rule $[14,15]$, which is the counterpart of isolated pentagon rule [16] of carbon fullerenes. Another class is of fullerene-like structures, based on combination of 5-and 6-membered rings, with the existence of $\mathrm{N}-\mathrm{N}$ and $\mathrm{B}-\mathrm{B}$ bonds.

[^0]According to our best knowledge, detailed bonding analysis is still infrequent in theoretical works on $(\mathrm{BN})_{\mathrm{n}}$ cages. On the other hand, the oligomers and solids of $(\mathrm{BN})_{\mathrm{n}}$ cages have attracted the interest of theorists [8]. Due to the enlarged-size of such systems, the cost of direct calculations is quite high and the possibility of predicting properties of oligomers just from studies on monomer is of great interest. So the purpose of this paper is, first, to analyse the details of bonding of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage [6, 8, 15], which is shown to be the most stable isomer [6]. Second, as a preliminary step to explore the possibility of forecasting properties of oligomers just from studies on monomer, we predict the relative stability of four dimers of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage, just from bonding analysis on single cage, and examine the validity of this prediction.

DFT [17] has become an efficient tool for theoretical studies on $(\mathrm{BN})_{\mathrm{n}}$ cage [3, 4, 6-11] and NBO analysis [1821] has been successfully applied into various molecular systems [22-34]. One thing worth noting is that conjugated systems, such as benzene and $\pi$-conjugated linear molecules, have been well studied with NBO analysis [32-34]. This fact points out the possibility of application of NBO analysis into alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage since this cage is the boron-nitride analogy of curved $\pi$-conjugated carbon fullerene. So bonding analysis in this paper is based on a combined study of DFT with NBO method. This combination has been shown to provide significant improvement over NBO analysis at HF level [20, 35, 36]. To ensure the reliability of such analysis, four DFT functionals, together with HF and MP2 [37] methods, are applied in this paper.

## Theoretical methods and computational details

NBO method transforms atomic orbitals into a new basis set of localized lewis(core, bonding) and non-lewis(mainly
antibonding) NBO orbitals [18-20]. In closed-shell systems, lewis NBOs are nearly double-occupied [18] and provide the most accurate natural Lewis structure(NLS) description of the system [19].

Since antibonding NBO is of non-zero occupancy and leads to certain energetic stabilization [18], it contributes to bonding interaction between bond-connected atoms as well as bonding NBO. Lower orbital energy(OE) means occupying electrons are more stable at zeroth-order approximation, thus stronger bonding NBO is related to lower OE. Since occupancies( Q ) of various NBOs are different, orbital energy multiplied with its occupancy(OEMO) is more suitable to measure the strength of bonding NBO. The strength of antibonding NBOs can be measured by deletion energy(DE) [21, 27] and second-order interaction energy (SOIE) [34] quantitatively. Higher value of DE or SOIE corresponds to stronger antibonding NBO.

Alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage is fully optimized by B3LYP [38, 39], BLYP [39, 40], MPW1PW91 and MPWPW91 [41, 42], together with HF and MP2 [37] method. 6-31G* basis set, which is proven to be suitable for $(\mathrm{BN})_{\mathrm{n}}$ cages $[9,10]$, is used in this paper. The local minimum character of optimized structures is confirmed by frequency analysis [43]. Gaussian 03 program [44] with NBO 5.0 code [45] is used for all calculations and analysis.

## Results and discussions

## The decomposion of bonding structure

The bonding structure of $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage could be decomposed into $\sigma$ skeleton, shown in Fig. 1, and $\pi$ bond system.

The $\sigma$ skeleton is considered as being constructed from $\mathrm{sp}^{2}$ hybrids and can be described by $\sigma$ bonding and $\sigma$ antibonding ( $\sigma^{*}$ ) NBOs. Similarly, the $\pi$ bond system, constructed from remaining $\pi$-type atomic orbitals, could be characterized in terms of $\pi$ and $\pi^{*}$ NBOs. Because of the cage curvature, $\pi$ bond system should be partially conjugated.

In order to simplify the bonding analysis, full usage of molecular symmetry is necessary. Due to the Td symmetry of the cage [15], confirmed by frequency analysis of optimized structures, both nitrogen and boron atoms can be separated into two types as shown in $\sigma$ skeleton in Fig. 1: type $i$-atoms lying on $\mathrm{C}_{3}$ axes(atom 25-atom 32 in slanted text) and type $j$-atoms constituting 4-membered rings(atom 1 -atom 24 in normal text). So there are only three types of BN units: $\mathrm{N}(i)-\mathrm{B}(j), \mathrm{B}(i)-\mathrm{N}(j)$ and $\mathrm{N}(j)-\mathrm{B}(j)$, noted as $\mathrm{BN}(1), \mathrm{BN}(2), \mathrm{BN}(3)$ respectively in this paper.

As shown in Table 1, the optimized bond lengths for alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage are shorter than BN single bond in $\mathrm{H}_{3} \mathrm{~B}-\mathrm{NH}_{3}$ but longer than double bond in $\mathrm{H}_{2} \mathrm{~B}=\mathrm{NH}_{2}$, supporting the existence of partially conjugated $\pi$ bond system. The order of bond lengths is: $\mathrm{BN}(3)>\mathrm{BN}(2)>$ BN(1).

As a matter of fact, an alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage with equal BN bonds is also obtained from point-wise optimization at B3LYP/6-31G* level. The calculations show that the energy of this structure is $200.91 \mathrm{kcal} / \mathrm{mol}$ higher than that of the minimum calculated here. Hence the monomeric $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage should be favorable to a distorted structure from a geometry with perfectly square $\mathrm{B}_{2} \mathrm{~N}_{2}$ and perfectly hexagonal $\mathrm{B}_{3} \mathrm{~N}_{3}$ faces. This is in accordance with the previous calculations for alternant $(\mathrm{BN})_{\mathrm{n}}$ cages with more than 20 atoms [4-6]. Moreover, 4-membered ring $\mathrm{B}_{2} \mathrm{~N}_{2}$ and


Fig. 1 Optimized geometry and $\sigma$ skeleton of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage

Table 1 Optimized bond lengths for alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage $^{\mathrm{a}}$

|  | B3LYP | BLYP | MPW1 | MPW | HF | MP2 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| BN(3) | 1.473 | 1.485 | 1.469 | 1.480 | 1.463 | 1.476 |
| BN(2) | 1.458 | 1.456 | 1.466 | 1.454 | 1.464 | 1.452 |
| BN(1) | 1.669 | 1.684 | 1.651 | 1.461 | 1.450 | 1.460 |
| B-N | 1.392 | 1.401 | 1.389 | 1.660 | 1.688 | 1.665 |
| $B=N^{\mathrm{c}}$ |  |  |  | 1.389 | 1.394 |  |

${ }^{\text {a }}$ Bond lengths in $\AA .{ }^{\mathrm{b}}$ length of BN single bond in $\mathrm{H}_{3} \mathrm{~B}-\mathrm{NH}_{3} .{ }^{\mathrm{c}}$ length of BN double bond in $\mathrm{H}_{2} \mathrm{~B}=\mathrm{NH}_{2}$

6-membered ring $\mathrm{B}_{3} \mathrm{~N}_{3}$ are also distorted by the same calculations. The deviation may be due to different geometry preference for B and N atoms which have different configurations of valent electron. Borons prefer planar geometry and nitrogens prefer pyramidalization [15]. However, for the 6 -membered ring $\mathrm{B}_{3} \mathrm{~N}_{3} \mathrm{H}_{6}$, a recent optimization leads to perfectly hexagonal structure [46]. As to the 4membered ring $\mathrm{B}_{2} \mathrm{~N}_{2} \mathrm{H}_{4}$, i.e., 1,3-diazadiboretidine, the optimized structures are also deviated from perfectly square faces [47].

By the way, experiments have shown the existence of boron-nitride analogy of graphite [48]. This kind of graphite-like boron nitride could be constructed from not only hexagonal but also rhombohedral rings [49, 50]. This reflects the diversity of boron nitride structures which is worth exploration.

Further decomposition into contributions from various NBOs

Selection of natural Lewis structure(NLS) Default NLS with different methods are not unique. Even nitrogen lone pairs and five-valency boron atoms exist in the default NLS of DFT calculations. For the consistency of consequent bonding analysis, a uniform NLS is necessary. Since there are three different types of BN unit and only 16 double bonds are available for 48 BN units of the $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage, a suitable NLS, without lone pairs and violations of lewis octet rule, is selected. In this NLS, one third of every distinct type of BN units are connected with double bonds as shown in Fig. 2.

Details of bonding interaction of $B N$ units For a certain BN unit, its bonding interaction is determined by all the NBOs within this unit as shown in Sect.2; thus it can be analyzed with the parameters of NBO as shown in Tables 2 and 3. Since only one third of BN units of one given type are connected with double bonds, shown by selected NLS (Fig. 2), $\pi$ and $\pi^{*}$ NBOs should be considered as orbitals into which $\pi$ bond system is condensed. In most cases, maximum relative deviations of parameters for NBOs of a certain type BN units are less than $2 \%$; thus this fact provides
the validity of selected NLS and only averaged values of these parameters are tabulated. From Tables 2 and 3, all the methods show consistent characters, though actual values are some different with different calculations.

Lower occupancy ( Q ) of bonding NBO means more electrons have been donated into antibonding NBOs [18], thus lower occupancy leads to stronger donor-ability. From Table 2, it can be seen that the values of occupancy are $1.910-1.972$ and $1.667-1.785$ for $\sigma$ and $\pi$ bonding NBOs respectively. Therefore, $\pi$ bonding NBOs possess higher donor-ability than $\sigma$ bonding NBOs. The order of occupancies of a given type of bonding NBOs with different methods is $\mathrm{HF} \geq$ Hybrid functionals $\geq \mathrm{GGA}$ functionals $\geq$ MP2. The lowest value of MP2 occupancies is due to the multi-determinant wavefunction in which the limiting value of occupancy is less than 2 for closed-shell system [18, 21, 51]. It is easy to understand the intermediate results provided by hybrid functional since it can be considered as combination of pure GGA and HF.

For the energetic parameters in Table 2, $\sigma-\mathrm{BN}(1)$ and $\sigma-$ $\mathrm{BN}(2)$ are of nearly the same magnitude for OE and OEMO. However this magnitude is significantly lower than that of $\sigma$ $\mathrm{BN}(3)$ by $\sim 50 \mathrm{kcal} / \mathrm{mol}$ and $\sim 100 \mathrm{kcal} / \mathrm{mol}$ for OE and


Fig. 2 Selected NLS for alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage

Table 2 Calculated parameters of bonding NBOs of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage

|  |  | B3LYP | BLYP | MPW1 | MPW | HF | MP2 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\sigma-\mathrm{BN}(1)$ | Q ${ }^{\text {a }}$ | 1.960 | 1.959 | 1.959 | 1.958 | 1.963 | 1.927 |
|  | OE ${ }^{\text {b }}$ | -433.55 | -387.75 | -477.68 | -396.09 | -569.13 |  |
|  | OEMO ${ }^{\text {c }}$ | -849.99 | -760.00 | -877.17 | -775.68 | -1117.16 |  |
| $\sigma-\mathrm{BN}(2)$ | Q | 1.972 | 1.972 | 1.971 | 1.971 | 1.972 | 1.938 |
|  | OE | -430.08 | -384.34 | -444.45 | -393.11 | -565.57 |  |
|  | OEMO | -847.99 | -757.81 | -876.02 | -774.81 | -1114.94 |  |
| $\sigma-\mathrm{BN}(3)$ | Q | 1.946 | 1.945 | 1.945 | 1.943 | 1.951 | 1.910 |
|  | OE | -381.87 | -338.43 | -394.98 | -346.39 | -510.46 |  |
|  | OEMO | -743.22 | -658.13 | -768.31 | -673.01 | -996.03 |  |
| $\pi$ - $\mathrm{BN}(1)$ | Q | 1.741 | 1.727 | 1.741 | 1.724 | 1.785 | 1.710 |
|  | OE | -211.51 | -183.22 | -220.74 | -190.70 | -288.15 |  |
|  | OEMO | -368.21 | -316.45 | -384.31 | -320.80 | -514.24 |  |
| $\pi$-BN(2) | Q | 1.701 | 1.688 | 1.697 | 1.680 | 1.747 | 1.667 |
|  | OE | -215.32 | -187.46 | -224.60 | -195.46 | -290.24 |  |
|  | OEMO | -366.27 | -316.42 | -381.15 | -328.37 | -507.03 |  |
| $\pi$ - $\mathrm{BN}(3)$ | Q | 1.709 | 1.694 | 1.706 | 1.688 | 1.758 | 1.678 |
|  | OE | -226.13 | -197.50 | -235.92 | -205.89 | -303.29 |  |
|  | OEMO | -386.45 | -334.60 | -402.58 | -347.50 | -533.16 |  |

${ }^{\text {a }}$ Occupancy in electron number ${ }^{\mathrm{b}}$ Orbital energy in $\mathrm{kcal} / \mathrm{mol}^{\mathrm{c}}$ Orbital energy multiplied with its occupancy in $\mathrm{kcal} / \mathrm{mol}$

OEMO respectively. The values of OE and OEMO of $\pi$-BN (1) are very close to those of $\pi-\mathrm{BN}(2)$ and higher than those of $\pi-\mathrm{BN}(3)$ by $\sim 13 \mathrm{kcal} / \mathrm{mol}$ and $\sim 20 \mathrm{kcal} / \mathrm{mol}$ for OE and OEMO respectively. For $\sigma$ bonding NBOs, the shorter the bond length is, the lower the OE and OEMO are. However it is just contrary for $\pi$ bonding NBOs. It seems that different types of bonding NBOs have different relation-
ship with bond lengths and this may be the result of curvature structure of this cage. Finally the order of strength of bonding NBOs is $\sigma-B N(1) \simeq \sigma-B N(2) \gg$ $\sigma-B N(3) \gg \pi-B N(3)>\pi-B N(1) \simeq \pi-B N(2)$.

As to the parameters of antibonding NBOs shown in Table 3, the occupancies of $\pi^{*}$ NBOs are much larger than those of $\sigma^{*}$ NBOs, i.e., $\pi^{*}$ NBOs have stronger acceptor-

Table 3 Calculated parameters of antibonding NBOs of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage

|  |  | B3LYP | BLYP | MPW1 | MPW | HF | MP2 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\sigma^{*}-\mathrm{BN}(1)$ | Q | 0.055 | 0.057 | 0.056 | 0.059 | 0.048 | 0.075 |
|  | DE ${ }^{\text {a }}$ | 38.99 | 36.51 | 40.95 | 38.15 | 45.53 |  |
|  | SOIE ${ }^{\text {b }}$ | 50.98 | 46.57 | 53.98 | 48.67 | 61.89 |  |
| $\sigma^{*}-\mathrm{BN}(2)$ | Q | 0.041 | 0.045 | 0.042 | 0.045 | 0.036 | 0.061 |
|  | DE | 27.10 | 25.32 | 28.61 | 26.45 | 32.52 |  |
|  | SOIE | 30.91 | 27.83 | 32.41 | 28.85 | 39.25 |  |
| $\sigma^{*}-\mathrm{BN}(3)$ | Q | 0.051 | 0.056 | 0.052 | 0.056 | 0.043 | 0.072 |
|  | DE | 31.12 | 26.74 | 29.92 | 28.14 | 32.82 |  |
|  | SOIE | 34.79 | 31.38 | 36.83 | 32.86 | 44.46 |  |
| $\pi^{*}$-BN(1) | Q | 0.244 | 0.256 | 0.244 | 0.261 | 0.200 | 0.261 |
|  | DE | 86.94 | 85.58 | 89.64 | 88.62 | 88.06 |  |
|  | SOIE | 128.50 | 117.22 | 133.31 | 119.98 | 157.39 |  |
| $\pi^{*}$-BN(2) | Q | 0.229 | 0.243 | 0.228 | 0.244 | 0.186 | 0.240 |
|  | DE | 67.28 | 66.28 | 68.69 | 67.58 | 69.36 |  |
|  | SOIE | 72.58 | 62.46 | 75.73 | 63.42 | 101.36 |  |
| $\pi^{*}-\mathrm{BN}(3)$ | Q | 0.233 | 0.244 | 0.233 | 0.246 | 0.196 | 0.249 |
|  | DE | 81.32 | 79.41 | 83.65 | 81.72 | 85.80 |  |
|  | SOIE | 101.62 | 89.58 | 106.34 | 92.17 | 137.44 |  |

[^1]ability compared to $\sigma^{*}$ NBOs. DE values demonstrate that the strength of $\pi^{*}-\mathrm{BN}(3)$ is close to that of $\pi^{*}-\mathrm{BN}(1)$ while markedly stronger than $\pi^{*}-\mathrm{BN}(2)$ by $15 \sim 20 \mathrm{kcal} / \mathrm{mol}$. $\sigma^{*}$ $\mathrm{BN}(3)$ and $\sigma^{*}-\mathrm{BN}(2)$ are of comparable strengths and weaker than $\sigma^{*}-\mathrm{BN}(1)$ by $\sim 10 \mathrm{kcal} / \mathrm{mol}$. Although SOIE values are higher than corresponding DE values, especially for $\pi^{*}$ NBOs, it is consistent with DE in the aspect of relative strength of different antibonding NBOs.

Summary of bonding analysis The order of strengths is $\sigma-\quad B N(1) \simeq \sigma-B N(2) \gg \sigma-B N(3) \gg \pi-B N(3)>$ $\pi-B N(1) \simeq \pi-B N(2)$ for bonding NBOs and $\pi *-$ $B N(1) \simeq \pi *-B N(3)>\pi *-B N(2) \gg \sigma-B N(1)>\sigma-$ $B N(3) \simeq \sigma-B N(2)$ for antibonding NBOs. So for $\sigma$ bonding interaction, it is $\mathrm{BN}(1)>\mathrm{BN}(2)>\mathrm{BN}(3)$ and for $\pi$ bonding interaction it is $\mathrm{BN}(3)>\mathrm{BN}(1)>\mathrm{BN}(2)$.

Bonding index analysis
In order to provide sufficient description on alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage, wiberg bond index (WBI) [21, 52] and overlap-weighted NAO bond order (OWBO) [21] are calculated and shown in Table 4.

The calculated values in Table 4 indicate that the B-N bonds of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage are just between the typical single and double BN bond, thus bond index also support that the $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage is a partially conjugated system. It should be pointed out that the order of bond index are not exactly matched with bond lengths. Especially, WBI of BN(3) with longer bond length is larger than that of $\mathrm{BN}(1)$ with shorter bond length. Although WBI has been used in the study of $(\mathrm{BN})_{10}$ [3], the results here suggest that bond indexes can only provide a primitive description of $(\mathrm{BN})_{\mathrm{n}}$ cage

Predicting relative stability of dimers
Based on the above bonding analysis we now discuss the possible stability for two cages joining to form covalent
dimers $\left(\mathrm{B}_{16} \mathrm{~N}_{16}\right)_{2}$. Four linking patterns, which are $\mathrm{BN}(1)_{a}=$ $\mathrm{BN}(1)_{b}, \mathrm{BN}(2)_{a}=\mathrm{BN}(2)_{b}, \mathrm{BN}(3)_{a}=\mathrm{BN}(3)_{b}$ and $\mathrm{BN}(1)_{a}=$ $\mathrm{BN}(2)_{b}$ with $a, b$ denoting different $\mathrm{B}_{16} \mathrm{~N}_{16}$ cages, are considered. These four dimers are abbreviated as D1, D2, D3, D4, shown in Fig. 3.

Assuming the reservation of $\sigma$ skeletons of monomer cages, the main change, after the formation of dimers, is that $\pi$ bonds of two BN units of original $\mathrm{B}_{16} \mathrm{~N}_{16}$ cages are broken and replaced by two inter-cage single bonds. If the strength of inter-cage bonds were almost the same, the relative stability of the dimers should be determined mainly by the relative strength of replaced $\pi$ bonds. Since the order of the strength of $\pi$ bond is $\mathrm{BN}(3)>\mathrm{BN}(1)>\mathrm{BN}(2)$ from bonding analysis on monomer, the relative stability of the $\left(\mathrm{B}_{16} \mathrm{~N}_{16}\right)_{2}$ dimers is predicted as $\mathrm{D} 2>\mathrm{D} 4>\mathrm{D} 1>\mathrm{D} 3$.

To verify this prediction, we optimized the four dimers at B3LYP/ 6-31G* level. The calculations give the stable order of $\mathrm{D} 3>\mathrm{D} 2>\mathrm{D} 4>\mathrm{D} 1$ as shown in Table 5. It can be seen that only D3 is not in accordance with the predicted order. This exception is because the $\sigma$ skeletons of monomer cages have been broken after the formation of D3. The lengths of original BN units, connected by intercage BN bonds, change to $2.586 \AA$ and is much longer than that of a usual single BN bond. Moreover, the inter-cage BN bonds are of nearly the same length as the double bond in $\mathrm{H}_{2} \mathrm{~B}=\mathrm{NH}_{2}(1.392 \AA)$ as shown in Table 5. Therefore in the case of D3, both $\sigma$ and $\pi$ bonding interactions of two original BN units are broken and replaced by two inter-cage double bonds. Although $\mathrm{BN}(3)$ is the strongest in the aspect of $\pi$ bonding interaction, its $\sigma$ bonding interaction is quite weaker than other types of BN units, considering the OEMO of $\sigma-\mathrm{BN}(3)$ is $\sim 100 \mathrm{kcal} / \mathrm{mol}$ higher than other two types of $\sigma$ NBOs (Table 2). This may be the reason why D3 becomes the most stable dimer.

The main difference between $\pi$ bond of $\mathrm{BN}(1)$ and BN (2) is that the contribution from $\pi^{*}-\mathrm{BN}(1)$ is stronger than that of $\pi^{*}-\mathrm{BN}(2)$ by $6.55 \mathrm{kcal} / \mathrm{mol}$ on average [53]. There are moderate deviations of calculated relative energies

Table 4 Calculated bond index for alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage

|  |  | B3LYP | BLYP | MPW1 | MPW | HF |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| BN(1) | WBI | 0.857 | 0.867 | 0.869 | 0.867 | 0.818 |
| BN(2) | OWBO | 0.900 | 0.902 | 0.905 | 0.904 | 0.882 |
|  | WBI | 0.875 | 0.886 | 0.889 | 0.888 | 0.828 |
| BN(3) | OWBO | 0.903 | 0.906 | 0.909 | 0.908 | 0.884 |
|  | WBI | 0.870 | 0.878 | 0.882 | 0.880 | 0.832 |
| $\mathrm{H}_{3} \mathrm{~B}-\mathrm{NH}_{3}{ }^{\text {a }}$ | OWBO | 0.884 | 0.884 | 0.888 | 0.814 |  |
| $\mathrm{H}_{2} \mathrm{~B}=\mathrm{NH}_{2}{ }^{\mathrm{b}}$ | WBI | 0.608 | 0.626 | 0.634 | 0.887 | 0.873 |
|  | OWBO | 0.594 | 0.610 | 0.615 | 0.619 | 0.559 |

[^2]Fig. 3 Optimized structures of the four $\left(\mathrm{B}_{16} \mathrm{~N}_{16}\right)_{2}$

from the estimation based on DE as shown in Table 5. This is because various inter-cage bonds are not exactly the same and monomers are deformed during the formation of dimer. From Table 5, D4 is less deformed than D2 and the strength of inter-cage bonds of D4 is stronger than that of D2, shown by bond lengths. Even though, D4 is still less stable than D2 and this fact stresses the possibility of forecasting properties of oligomers just from studies on monomer.

As a comparison, we also explored the dimerization of 1,3-diazadiboretidine. Here, only one isomer of $\mathrm{B}_{4} \mathrm{~N}_{4} \mathrm{H}_{8}$, which is a cage structure with alternant BN units [54], was considered. Our calculations show that, this dimerization is exothermic, about $24.49 \mathrm{kcal} / \mathrm{mol}$, which is consistent with MP2 calculations [54]. The dimerization of alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage is also exothermic, $11.92-55.88 \mathrm{kcal} / \mathrm{mol}$ for the dimers studied here. Previous work reported the possibility of delocalization of electrons in 1,3-diazadiboretidine [47], which may result into the partially conjugated $\pi$ bond system similar to alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage. However, both boron and nitrogen atoms in the cage structure of $\mathrm{B}_{4} \mathrm{~N}_{4} \mathrm{H}_{8}$ are 4-coordinated, thus this structure excludes the possibility
for conjugated $\pi$ bond. This is different from the dimers of $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage.

## Conclusions

Based on a combined study of DFT with NBO method, a detailed bonding analysis on alternant $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage is performed. First the bonding structure is seperated into two components: $\sigma$ skeleton and $\pi$ bond system. Then bonding interaction of one single BN unit is decomposed into contributions from various NBOs. Results with different methods agree with each other well and details of bonding interactions of any BN unit are obtained in this way.

Relative stability of four covalent dimers of the $\mathrm{B}_{16} \mathrm{~N}_{16}$ cage is predicted from bonding analysis on monomers. Moreover, direct calculations on dimers verify this prediction if the $\sigma$ skeletons of monomer cages are not broken. Hence the possibility of forecasting properties of oligomers just from studies on the monomer, which is a more efficient way compared to direct calculations on oligomers, is highlighted by this study.

Table 5 Relative energies and bond lengths of the four dimers ${ }^{a}$

|  | D1 | D2 | D3 | D4 |
| :--- | :--- | :--- | :--- | :--- |
| Relative energy | $0.00(0.00)$ | $-8.96(-13.10)$ | -43.96 | $-5.45(-6.55)$ |
| Inter-cage BN | $1.613 / 1.613$ | $1.596 / 1.596$ | $1.404 / 1.404$ | $1.581 / 1.581$ |
| Deformation energy $^{\text {b }}$ | 51.26 | 56.02 |  | 53.79 |

[^3]Bond index analysis is also carried out with the conclusion that bond indexes can only provide a primitive description of boron nitride cages.

Acknowledgements We gratefully acknowledge Professor Frank A. Weinhold (University of Wisconsin-Madison, Theoretical Chemistry Institute) for his kindly answers to our quetions on combined study of DFT with NBO method. This work is supported by the Research Fund for the Doctoral Program of Higher Education (20060027001) and the Major State Basic Research Development Programs (Grant No. 2002CB613406).

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[^0]:    B. Yin $\cdot G$. Wang $\cdot N$. Sa $\cdot$ Y. Huang $(\boxtimes)$

    College of Chemistry, Beijing Normal University, Beijing 100875, People's Republic China
    e-mail: yuanhe@bnu.edu.cn

[^1]:    ${ }^{\text {a }}$ Deletion energy in kcal/mol ${ }^{\mathrm{b}}$ Second-order interaction energy in kcal/mol

[^2]:    ${ }^{\mathrm{a}}$ Bond index for BN unit of $\mathrm{H}_{3} \mathrm{~B}-\mathrm{NH}_{3}{ }^{\mathrm{b}}$ Bond index for BN unit of $\mathrm{H}_{2} \mathrm{~B}=\mathrm{NH}_{2}$

[^3]:    ${ }^{\text {a }}$ Energy in $\mathrm{kcal} / \mathrm{mol}$ and bond length in $\AA$. Relative energies, based on DE estimation exclusively, are shown in parenthesis. ${ }^{\mathrm{b}}$ The difference between energy of two monomers with structures fixed as in dimers and that of two monomers with optimal structures

